

1.

a)

$$\vec{e}_1 = (0, 1, 0) \quad \vec{e}_2 = (\cos \alpha, 0, -\sin \alpha)$$

Note that signs are not fixed by the specifications of this problem.

b) Normal dispersion

$$\theta_c \approx 44 \quad \theta_B \approx 56$$

in degrees

c) Many possible answers, e.g.: sound wave in gas is longitudinal, electromagnetic is transverse,...

d) E.g. plane waves and spherical waves, also bessel beams. BUT: Not cylindrical waves in the form given in the class — because those are approximations.

Note that this is poorly worded question: You could give sol_1 as an exact solution to wave, and claim $1.1sol_1$ is another exact solution — but I was obviously not asking that.

e) It is useful to remember that one micron corresponds to roughly $E_{ph} = 1.24 \text{ eV}$ which approximately $1.99 \times 10^{-19} \text{ J}$.

Photon flux: $P/E_{ph} = 0.5 \times 10^{19} \text{ s}^{-1}$.

f) Angular momentum is transferred to the plate.

g) Airy disk is the bright central spot of the far-field diffraction pattern for a circular apperture. Equivalently, it may be the bright spot in the focus of a lens (with an apperture).

h)

One of the slits closed: $I = I_b/4$.

Slits illuminated by perpendicularly polarized light: $I = I_b/2$.

2.

b) Use divergence equation (Gauss laws). Calculation:

$$0 = \nabla \cdot \vec{E} = \nabla \cdot \hat{e} E_0 e^{i\vec{k} \cdot \vec{r} - i\omega t} = i\vec{k} \cdot \hat{e} E_0 e^{i\vec{k} \cdot \vec{r} - i\omega t}$$

which implies:

$$\vec{k} \cdot \hat{e} = 0$$

I.E. the wave-vector and the polarization vector are perpendicular.

c)

$$\nabla \times \vec{H} = \partial_t \vec{D}$$

$$i\vec{k} \times \vec{H} = -i\omega \vec{D}$$

$$i\vec{k} \times \vec{B} = -i\omega \mu_0 \epsilon_0 \epsilon_r \vec{E}$$

$$i\vec{k} \times \hat{b} B_0 = -i\omega \mu_0 \epsilon_0 \epsilon_r \hat{e} E_0$$

$$i\hat{k} \times \hat{b} |k| B_0 = -i\omega \mu_0 \epsilon_0 \epsilon_r \hat{e} E_0$$

separate the vector and amplitude portions of the above to get:

$$\hat{k} \times \hat{b} = -\hat{e} \quad |k| B_0 = \omega n(\omega)/c B_0 = \omega \mu_0 \epsilon_0 \epsilon_r E_0 = \omega/c^2 n(\omega)^2 E_0 \rightarrow c/n(\omega) B_0 = E_0$$

3:

- a) Insert $\theta = 0$. NOTE: that there may be a different convention in which they will differ in the sign. It does NOT mean physical difference, of course.
- b) Evaluating:

$$r_s = \frac{1-n}{1+n} \rightarrow r_s = \frac{1-1/n}{1+1/n} = -\frac{1-n}{1+n}$$

so the sign changes, but intensity reflectivity scales with r^2 so it has no effect.

- d) The jones matrix is diagonal with reflection coefficients on diagonal, because polarization s or p is preserved on reflection. Note that signs may depend on the orientation of the chosen frame of reference.
- e) Use Brewster angle of incidence, p polarization.

4:

- a) It is possible to have $T = 1$ when exactly at resonance $\Delta/2 = k\pi$. Finite reflectivity does not contradict, because there is infinite number of reflections from mirrors, each contributing to the transmission.
- b)

$$\mathcal{F} = \pi \frac{\sqrt{R}}{1-R} \quad \frac{c}{2d} \quad \frac{2d}{\lambda}$$

5:

- e) It is the second picture. The reason is that each set of edges contributes a perpendicular streak to the diffraction pattern. In the second picture, the diagonal streak is perpendicular to the hypotenuse of the triangle aperture.
- f) The symmetry that MUST be present in the Fraunhofer pattern is that if a point X, Y has intensity I , the same intensity must be found at point $-X, -Y$. This is called inversion symmetry. The pattern shown does not have this symmetry, it is therefore NOT a Fraunhofer diffraction pattern.