

Example: Extraordinary wave index, dependence on direction of propagation

This problem example shows calculation of the refractive index experienced by a plane wave propagating in a general direction through a uni-axial crystal.

Consider a uni-axial crystal with ordinary and extraordinary refractive indices n_o and n_e . A plane wave is propagating at an angle α with respect to the optic axis. Assume that the optic axis points along the z -axis of the coordinate system, and that the propagation vector lies in the $y - z$ plane.

Your task is to demonstrate that the refractive index depends on α as

$$n(\alpha) = \frac{n_o n_e}{\sqrt{n_e^2 \cos^2 \alpha + n_o^2 \sin^2 \alpha}}$$

Hint:

You will need to use all Maxwell equations, plus the medium constitutive relation,

$$\begin{aligned} D_x &= \epsilon_0 \epsilon_x E_x & D_y &= \epsilon_0 \epsilon_y E_y & D_z &= \epsilon_0 \epsilon_z E_z \\ \epsilon_x &= \epsilon_y = n_o^2 & \epsilon_z &= n_e^2 \end{aligned}$$

Lesson: The following steps show that the refractive index is obtained from the Maxwell as a solvability condition in the form of the angle-dependent dispersion relation:

$$k^2 = \frac{\omega^2 n(\alpha)^2}{c^2}$$

A) What is the unit vector \hat{n} in the direction of propagation?

Answer:

$$\hat{n} = (0, \sin \alpha, \cos \alpha)$$

B) Assuming that the wave is ordinary, determine the unit polarization direction vector \hat{o}

Answer:

$$\hat{o} = (1, 0, 0);$$

C) In the rest of the problem we assume that the wave is polarized extraordinary. Determine the polarization direction vector \hat{d} along which \vec{D} oscillates.

Answer:

We know that $\nabla \cdot \vec{D} = 0$ which implies that $\hat{d} \cdot \hat{n} = 0$. Moreover, extraordinary wave has polarization perpendicular to that of ordinary, which means that $d_x = 0$. These two properties imply (up to a sign):

$$\hat{d} = (0, -\cos \alpha, +\sin \alpha)$$

D) Let the amplitude of \vec{D} is D_0 . Denote ω and k the angular frequency and the magnitude of the propagation vector. Write down the expression for the non-zero components of \vec{D} .

Answer:

$$D_y = -D_0 \cos \alpha \exp[-i\omega t + k\hat{n} \cdot \vec{r}] \quad D_z = +D_0 \sin \alpha \exp[-i\omega t + k\hat{n} \cdot \vec{r}]$$

E) Write down non-zero components of the electric field

Answer:

$$E_y = -D_0/(\epsilon_0 n_o^2) \cos \alpha \exp[-i\omega t + k\hat{n} \cdot \vec{r}] \quad E_z = +D_0/(\epsilon_0 n_e^2) \sin \alpha \exp[-i\omega t + k\hat{n} \cdot \vec{r}]$$

F) Using Faraday, calculate the nonzero components of \vec{B}

Answer:

$$-\partial_t B_x = \partial_y E_z - \partial_z E_y$$

$$i\omega B_x = ik\hat{n}_y E_z - ik\hat{n}_z E_y = ik \sin \alpha E_z - ik \cos \alpha E_y$$

$$B_x = \frac{kD_0}{\epsilon_0 \omega} \left[\frac{\cos^2 \alpha}{n_o^2} + \frac{\sin^2 \alpha}{n_e^2} \right] \exp[-i\omega t + k\hat{n} \cdot \vec{r}]$$

G) Next, insert previous results in the Ampere. You should obtain two equivalent equations. The y component, for example, gives the following:

Answer:

$$\begin{aligned}\partial_t D_y &= \frac{1}{\mu_0} \partial_z B_x \\ -i\omega D_y &= \frac{1}{\mu_0} i k \hat{n}_z B_x \\ +D_0 \cos \alpha \exp[-i\omega t + k \hat{n} \cdot \vec{r}] &= \cos \alpha \frac{k^2 D_0}{\epsilon_0 \mu_0 \omega^2} \left[\frac{\cos^2 \alpha}{n_o^2} + \frac{\sin^2 \alpha}{n_e^2} \right] \exp[-i\omega t + k \hat{n} \cdot \vec{r}] \\ 1 &= \frac{k^2 c^2}{\omega^2} \left[\frac{\cos^2 \alpha}{n_o^2} + \frac{\sin^2 \alpha}{n_e^2} \right]\end{aligned}$$

H) Finally, parametrizing propagation number k in terms of the (angle-dependent) refractive index $n(\alpha)$ as

$$k^2 = n(\alpha)^2 \omega^2 / c^2 ,$$

obtain the dependence on the angle of propagation:

$$n(\alpha) = \frac{n_o n_e}{\sqrt{n_e^2 \cos^2 \alpha + n_o^2 \sin^2 \alpha}}$$