

OPTI-310. Review of useful formulas II.

Light-matter interactions

$$\vec{D} = \epsilon \vec{E} = \epsilon_0 \epsilon_r \vec{E} = \epsilon_0 n(\omega)^2 \vec{E} = \epsilon_0 \vec{E} + \vec{P}$$

Polarization (dipole moment density):

$$\vec{P} = N \vec{d} = N q_e \vec{x}(t)$$

Single oscillator model:

$$m_e \frac{d^2 \vec{x}}{dt^2} = q_e \vec{E}(t) - m_e \omega_0^2 \vec{x}(t)$$

$$n(\omega)^2 = 1 + \frac{N q_e^2}{\epsilon_0 m_e} \frac{1}{\omega_0^2 - \omega^2}$$

relation to susceptibility

$$n(\omega)^2 = 1 + \chi(\omega)$$

Resonance frequency:

$$\hbar \omega_0 = E_e - E_g$$

Generalization 1: Multiple oscillators, oscillator strengths

$$n(\omega)^2 = 1 + \frac{N q_e^2}{\epsilon_0 m_e} \sum_j \frac{f_j}{\omega_{0j}^2 - \omega^2}$$

Generalization 2: Including damping

$$n(\omega)^2 = 1 + \frac{N q_e^2}{\epsilon_0 m_e} \sum_j \frac{f_j}{\omega_{0j}^2 - \omega^2 + i\omega\gamma_j}$$

Generalization 3: Condensed medium (oscillators experience each other's field):

$$\frac{n(\omega)^2 - 1}{n(\omega)^2 + 2} = \frac{N q_e^2}{3\epsilon_0 m_e} \sum_j \frac{f_j}{\omega_{0j}^2 - \omega^2 + i\omega\gamma_j}$$

Generalization 4: Anisotropic oscillators for birefringent media

$$n_x(\omega)^2 = 1 + \frac{N q_e^2}{\epsilon_0 m_e} \frac{1}{\omega_{0x}^2 - \omega^2}$$

Cubic, isotropic

$$n_x = n_y = n_z$$

Uniaxial

$$n_x = n_y \neq n_z$$

Biaxial

$$n_x \neq n_y \neq n_z \neq n_x$$

Cauchy formula:

$$n(\omega) = C_1 + \frac{C_2}{\lambda^2} + \frac{C_3}{\lambda^4} + \dots$$

Sellmeier formula:

$$n^2(\omega) = 1 + \sum_j \frac{A_j \lambda^2}{\lambda^2 - \lambda_{0j}^2}$$

Helmholtz equation

$$\vec{E}(\vec{r}, t) = \vec{E}(\vec{r}, \omega) e^{-i\omega t}$$

$$[\nabla^2 + k^2] \vec{E}(\vec{r}, \omega) = 0$$

Boundary conditions at material interfaces

Tangential components of intensities \vec{E} , \vec{H} are continuous.

Normal components of inductions \vec{D} , \vec{B} are continuous.

Reflection and refraction at a material interface

$$n_i \sin \Theta_i = n_t \sin \Theta_t \quad , \quad \Theta_i = \Theta_r$$

Fresnel formulae

(Note: vector orientation convention as in Hecht)

\vec{E} normal to the plane of incidence, TE or s-polarization:

$$r_{\perp} = \frac{n_i \cos \Theta_i - n_t \cos \Theta_t}{n_i \cos \Theta_i + n_t \cos \Theta_t} \quad , \quad t_{\perp} = \frac{2n_i \cos \Theta_i}{n_i \cos \Theta_i + n_t \cos \Theta_t}$$

\vec{E} in the plane of incidence, TM or p-polarization:

$$r_{\parallel} = \frac{n_t \cos \Theta_i - n_i \cos \Theta_t}{n_t \cos \Theta_i + n_i \cos \Theta_t} \quad , \quad t_{\parallel} = \frac{2n_i \cos \Theta_i}{n_t \cos \Theta_i + n_i \cos \Theta_t}$$

Reflectance and transmittance, material and footprint factor

$$R = |r|^2 \quad , \quad T = \frac{n_t \cos \Theta_t}{n_i \cos \Theta_i} |t|^2$$

Brewster angle:

$$\tan \Theta_{Brewster} = n \equiv n_t / n_i$$

Critical angle:

$$\sin \Theta_{Critical} = n$$

Other forms:

$$r_{\perp} = -\frac{\sin(\Theta_i - \Theta_t)}{\sin(\Theta_i + \Theta_t)} \quad , \quad t_{\perp} = +\frac{2 \sin \Theta_t \cos \Theta_i}{\sin(\Theta_i + \Theta_t)}$$

$$r_{\parallel} = +\frac{\tan(\Theta_i - \Theta_t)}{\tan(\Theta_i + \Theta_t)} \quad , \quad t_{\parallel} = +\frac{2 \sin \Theta_t \cos \Theta_i}{\sin(\Theta_i + \Theta_t) \cos(\Theta_i - \Theta_t)}$$

$$r_s = \frac{\cos \Theta_i - \sqrt{n^2 - \sin^2 \Theta_i}}{\cos \Theta_i + \sqrt{n^2 - \sin^2 \Theta_i}} \quad , \quad t_s = \frac{2 \cos \Theta_i}{\cos \Theta_i + \sqrt{n^2 - \sin^2 \Theta_i}}$$

$$r_p = \frac{n^2 \cos \Theta_i - \sqrt{n^2 - \sin^2 \Theta_i}}{n^2 \cos \Theta_i + \sqrt{n^2 - \sin^2 \Theta_i}} \quad , \quad t_p = \frac{2n \cos \Theta_i}{\cos \Theta_i + \sqrt{n^2 - \sin^2 \Theta_i}}$$

Normal incidence:

$$R = \left(\frac{n_t - n_i}{n_t + n_i} \right)^2 \quad , \quad T = \frac{4n_t n_i}{(n_t + n_i)^2} \quad , \quad R + T = 1$$

Conversion: Hecht \rightarrow Fowles

$$r_{\perp} \quad , \quad t_{\perp} \rightarrow r_s \quad , \quad t_s$$

$$r_{\parallel} \quad , \quad t_{\parallel} \rightarrow -r_p \quad , \quad t_p$$

Total internal reflection

$$r_s = \frac{\cos \Theta_i - i \sqrt{\sin^2 \Theta_i - n^2}}{\cos \Theta_i + i \sqrt{\sin^2 \Theta_i - n^2}} \quad , \quad \Theta_i > \Theta_{cr} \quad , \quad |r_s| = 1$$

Polarization

$$E_x(z, t) = A_x \cos[kz - \omega t + \phi_x] \quad , \quad E_y(z, t) = A_y \cos[kz - \omega t + \phi_y]$$

Implicit equation: Ellipse for the “tip” of the electric vector

$$\left(\frac{E_x}{A_x} \right)^2 + \left(\frac{E_y}{A_y} \right)^2 - 2 \left(\frac{E_x}{A_x} \right) \left(\frac{E_y}{A_y} \right) \cos \phi_{yx} = \sin^2 \phi_{yx} \quad \phi_{yx} = \phi_y - \phi_x$$

Ellipse angle

$$\tan 2\alpha = \frac{2A_x A_y \cos \phi_{yx}}{A_x^2 - A_y^2}$$

RHC,LHC

$$\vec{E} = A \hat{e}_{\pm} e^{ikz - \omega t} + c.c.$$

$$\hat{e}_{\pm} \frac{1}{\sqrt{2}} (\hat{x} \mp i\hat{y})$$

Jones vectors
general:

$$\mathcal{E} = \begin{pmatrix} E_x \\ E_y \end{pmatrix}$$

linear - x

$$\begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

linear - y

$$\begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

RHC

$$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -i \end{pmatrix}$$

LHC

$$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ +i \end{pmatrix}$$

Jones matrices

Sequence of optical elements:

$$M = M_n \dots M_3 M_2 M_1$$

$$\mathcal{E}_{out} = M \mathcal{E}_{in}$$

Polarizer - x

$$\begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

Polarizer - y

$$\begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

Polarizer - ± 45

$$\frac{1}{2} \begin{pmatrix} 1 & \pm 1 \\ \pm 1 & 1 \end{pmatrix}$$

Quarter waveplate, fast axis vertical, horizontal, ± 45

$$\begin{pmatrix} 1 & 0 \\ 0 & -i \end{pmatrix} \quad , \quad \begin{pmatrix} 1 & 0 \\ 0 & +i \end{pmatrix} \quad , \quad \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & \pm i \\ \pm i & 1 \end{pmatrix}$$

Polarizer circular right, left

$$\frac{1}{2} \begin{pmatrix} 1 & +i \\ -i & 1 \end{pmatrix} \quad , \quad \frac{1}{2} \begin{pmatrix} 1 & -i \\ +i & 1 \end{pmatrix}$$

Rotator, rotating any (polarization-state) vector in positive, anti-clock-wise sense

$$\begin{pmatrix} \cos \alpha & -\sin \alpha \\ +\sin \alpha & \cos \alpha \end{pmatrix}$$

Rotated elements

$$R(\alpha) = \begin{pmatrix} \cos \alpha & +\sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix}$$

$$M = R^{-1}(\alpha)M'R(\alpha) = R(-\alpha)M'R(\alpha)$$

Rotated x-polarizer

$$M' = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}, \quad M = \begin{pmatrix} \cos^2 \alpha & \cos \alpha \sin \alpha \\ \cos \alpha \sin \alpha & \sin^2 \alpha \end{pmatrix}$$

Energy conservation and determinants, polarizers vs plates/retarders:

$$\det(M) = 0 \quad \text{vs} \quad |\det(M)| = 1$$